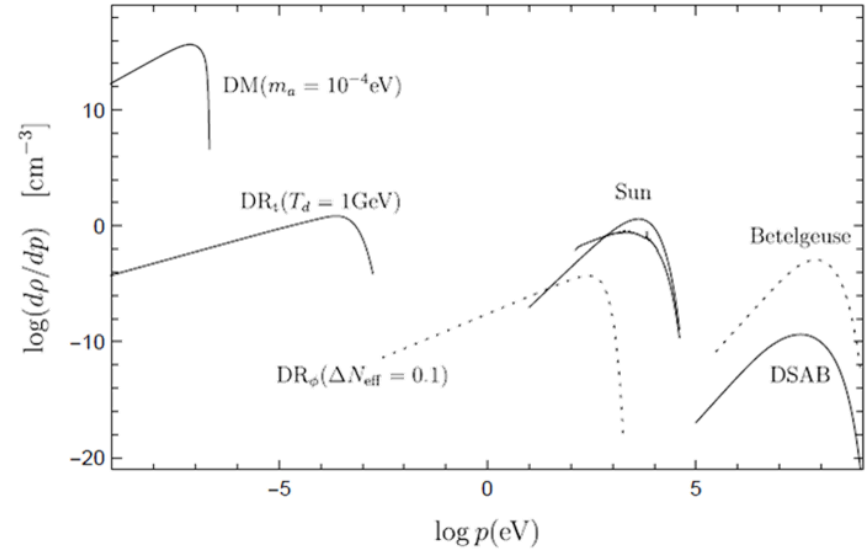


# AXION実験 CARRACK

小川 泉

# ALPを探る

- 何を？
  - 天然系
    - Solar axion
    - **Halo axion**
    - (Supernova axion...)
  - 作ってまえ
    - 実験室で作って検出



arXiv: 1801.08127

- どうやって？
  - 通常物質との結合

- 光子  $g_{a\gamma\gamma} = \frac{\alpha/2\pi}{f_{PQ}/N} (E_{PQ}/N - 1.95)$

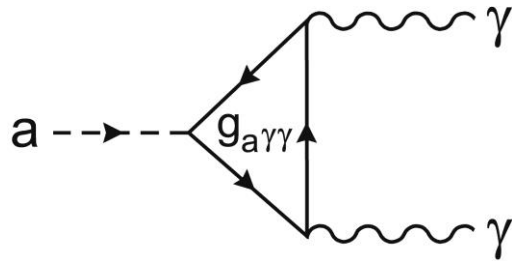
- 電子  $g_{aee} = \left( \frac{X_e}{N} + \frac{3\alpha^2}{4\pi} \left( \frac{E_{PQ}}{N} \ln \left( \frac{f_{PQ}}{m_a} \right) - 1.95 \ln \left( \frac{\Lambda_{\text{QCD}}}{m_e} \right) \right) \right) \left( \frac{m_e}{f_{PQ}/N} \right)$

- 陽子  $g_{app} = \dots$

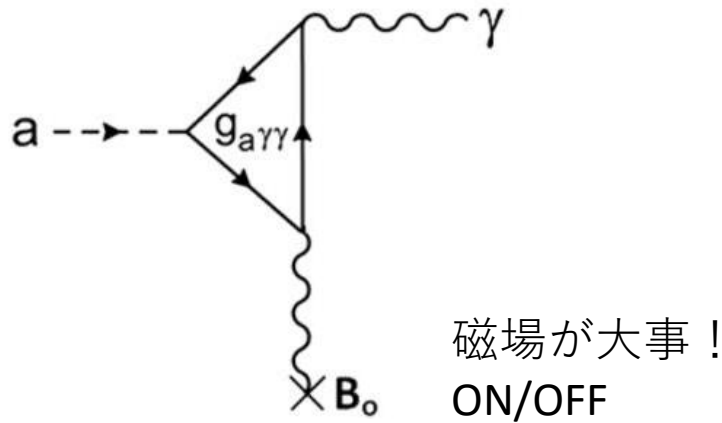
- 中性子  $g_{ann} = \dots$

# 四天王の中で最強： $g_{a\gamma\gamma}$

※最強：使い勝手がいい



- 実験では... (逆) プリマコフ効果
  - Helioscope, Haloscope,...



磁場が大事！  
ON/OFF

+resonance, boost....

# 現状

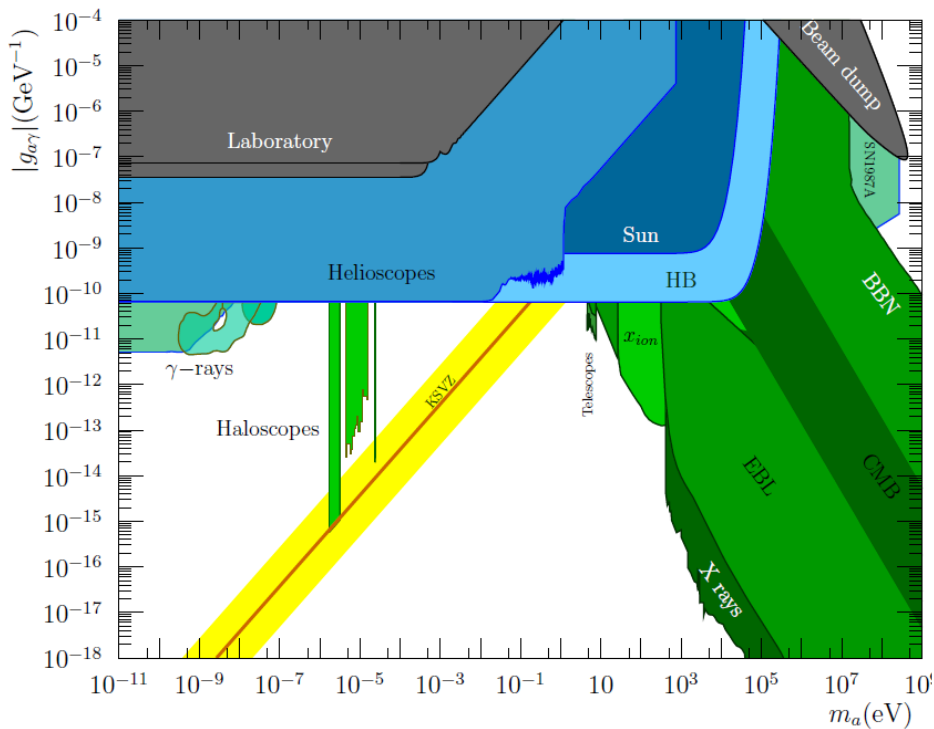
arXiv: 1801.08127

Detection method	$g_{a\gamma}$	$g_{ae}$	$g_{aN}$	$g_{A\gamma n}$	$g_{a\gamma}g_{ae}$	$g_{a\gamma}g_{aN}$	$g_{ae}g_{aN}$	$g_N\bar{g}_N$	Model dependency
Light shining through wall	×								no
Polarization experiments	×								no
Spin-dependent 5th force			×				×	×	no
Helioscopes	×				×	×			Sun
Primakoff-Bragg in crystals	×				×				Sun
Underground ion. detectors	×	×	×			×	×		Sun*
Haloscopes	×								DM
Pick up coil & LC circuit	×								DM
Dish antenna & dielectric	×								DM
DM-induced EDM (NMR)			×	×					DM
Spin precession in cavity		×							DM
Atomic transitions		×	×						DM

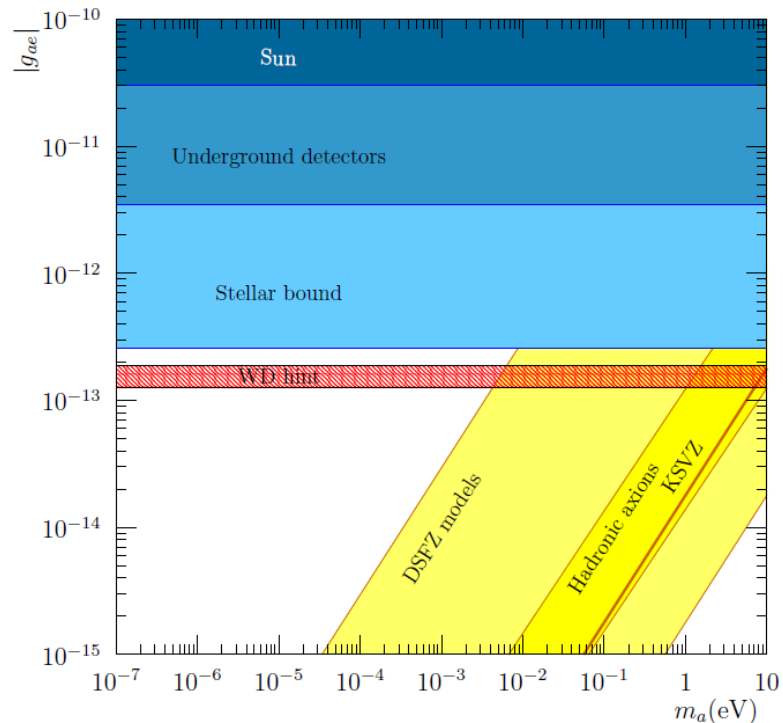
Table 3: List of the axion detection methods discussed in the review, with indication of the axion couplings (or product of couplings) that they are sensitive to, as well as whether they rely on astrophysical (axions/ALPs are produced by the Sun) or cosmological (the dark matter is made of axions/ALPs) assumptions. \*Also “DM” when searching for ALP DM signals, see section 6.2

# 現状

axion-photon coupling



axion-electron coupling



arXiv: 1801.08127

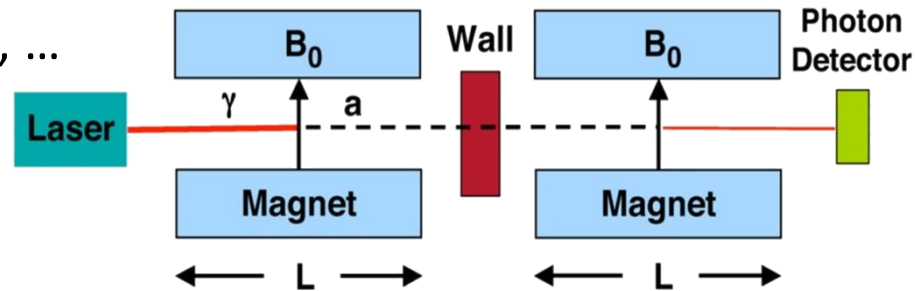
- 黒系：実験室
- 青系：天体物理（太陽、恒星他）
- 緑系：宇宙物理（DM halo, その他）
- 黄系：理論モデル（axion）

# 実験室での探索

- 利点：モデル（天体／宇宙物理）に依存しない
- 欠点：（現時点では）感度が低い
- 実験手法

- Light-shining-through wall (LSW)実験

- ALPS, CROWS, OSQAR, STAX, ...
- Laser, マイクロ波, X線



- Polarization実験

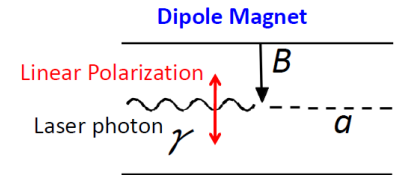
- PVLAS, ...

(1) Polarization Measurement

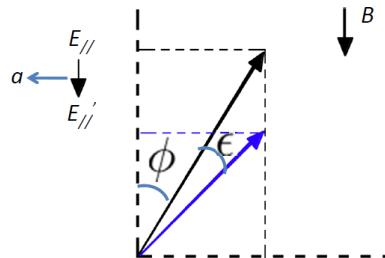
PVLAS, Q&A, ...

$$\vec{E} // \vec{B} \rightarrow \vec{E} \cdot \vec{B} \rightarrow \text{Pseudo-scalar}$$

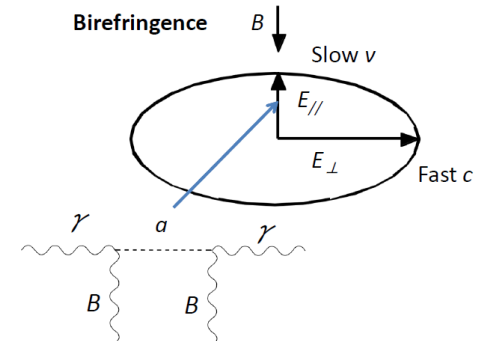
$$\vec{E} \perp \vec{B} \rightarrow \vec{E}^2 - \vec{B}^2 \rightarrow \text{Scalar}$$

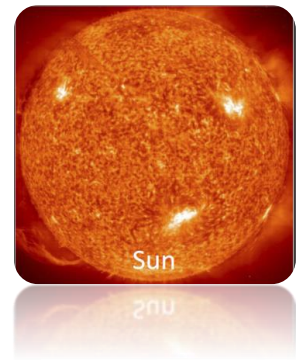


Polarization Rotation



Birefringence





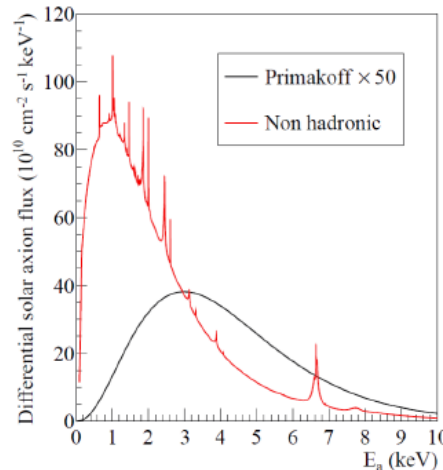
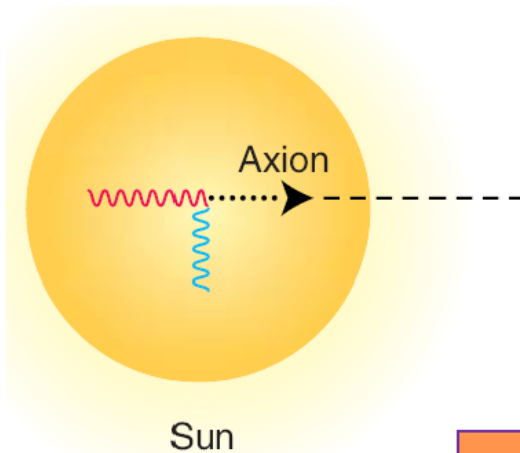
# 太陽axionの探索

$$\mathcal{P}(a \rightarrow \gamma) = 2.6 \times 10^{-17} \left( \frac{g_{a\gamma}}{10^{-10} \text{GeV}^{-1}} \right)^2 \left( \frac{B_e}{10 \text{ T}} \right)^2 \left( \frac{L}{10 \text{ m}} \right)^2 \mathcal{F}(qL)$$

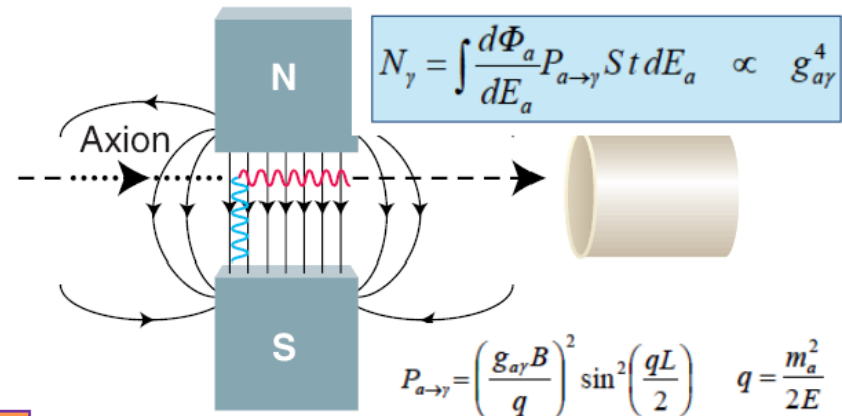
## ■ First axion helioscope proposed by P. Sikivie

Sikivie *PRL* 51:1415 (1983)

- Blackbody photons (keV) in solar core can be converted into axions in the presence of strong electromagnetic fields in the plasma
- Reconversions of axions into x-ray photons possible in strong laboratory magnetic field

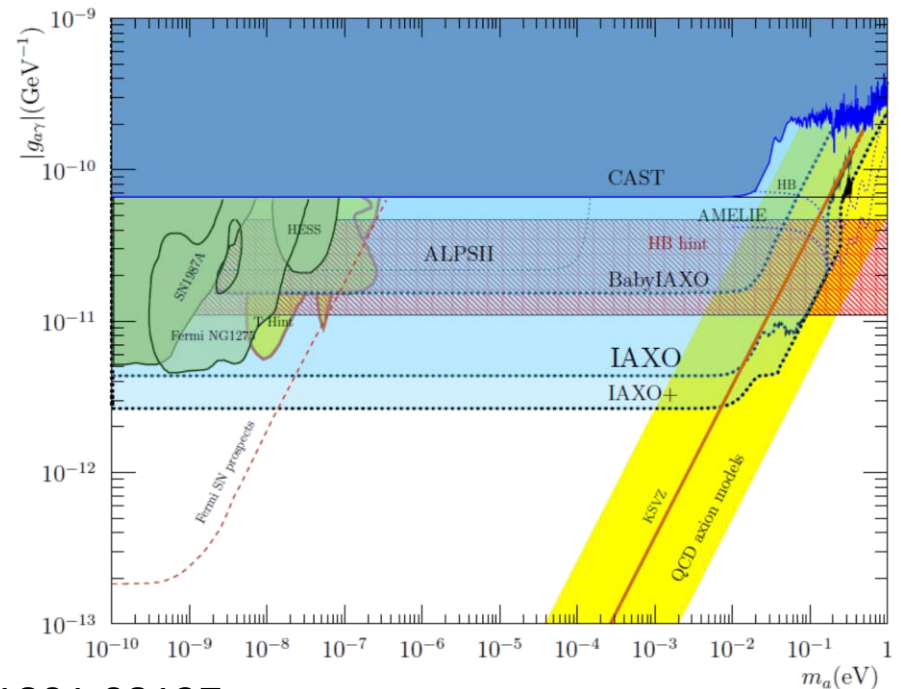
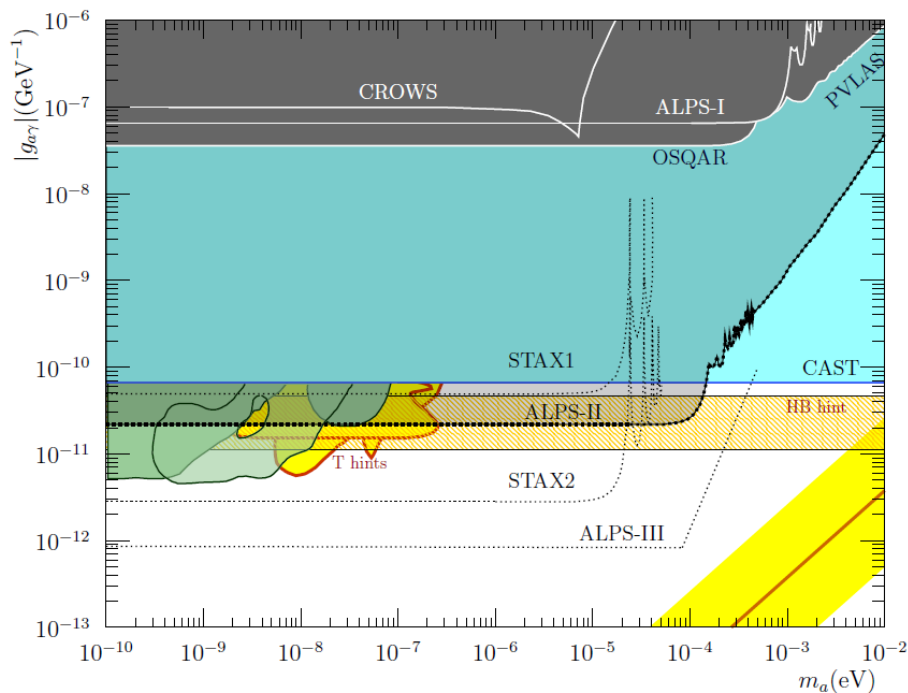


Redondo *JCAP* 1312 008 (2013)



# DH axion以外の 計画まとめ

	Experiment	status	B(T)	L(m)	Power(W)	A(cm <sup>2</sup> )
Laboratory	ALPS-I	Done	5	4.3	4	
	CROWS	Done	3	0.15	50	
	OSQAR	Ongoing	9	14.3	18.5	
	ALPS-II	In preparation	5	100	30	
	STAX	Concept	15	0.5	10 <sup>5</sup>	
Solar	Brookhaven	done	2.2	1.8		130
	SUMICO	done	4	2.5		18
	CAST	Ongoing	9	9.3		30
	BabyIAXO	In design	~2.5	10		2.8x10 <sup>3</sup>
	IAXO	In design	~2.5	22		2.8x10 <sup>4</sup>





# ダークハロー アクション

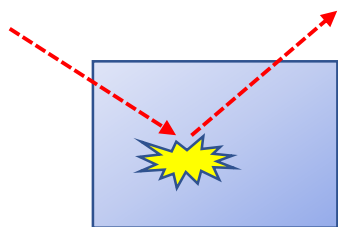
- 太陽近傍

$$\rho_{DM} = 0.2 - 0.56 (0.45) \text{ GeV/cm}^3$$

$$\beta \sim 10^{-3}$$

$m_{DM}$   
1 eV

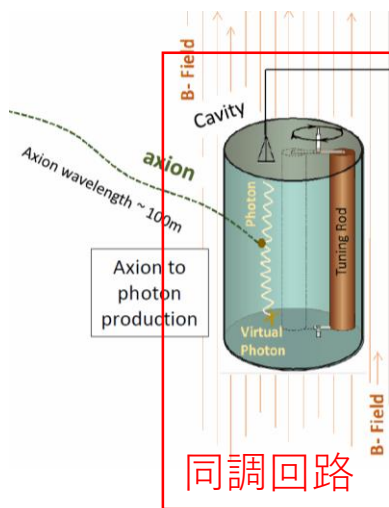
暗黒粒子  
WIMPs



反跳信号

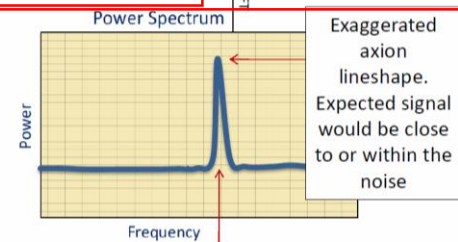
ド・ブロイ波長  
~1mm

暗黒波  
ALPs



検波・増幅回路

Amplify Digitize

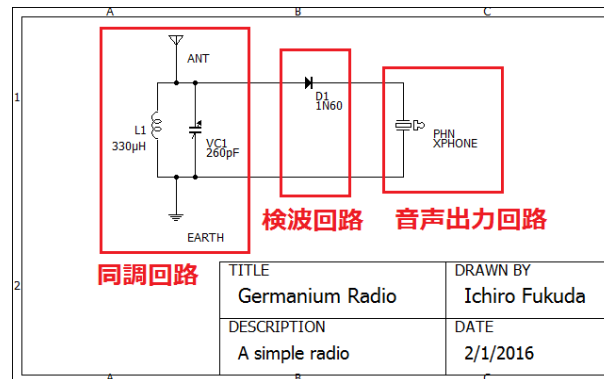


Axion mass unknown requiring a tunable resonator

Image Source: C. Boutan Thesis(2017)

記録・解析

同調回路

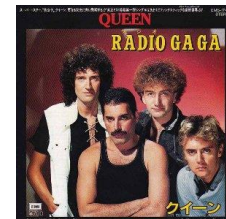


ゲルマニウムラジオ

<http://ele-tech.net/germanium-radio2/>

N.Woolet, talk at Rencontres de Moriond (2018)

Now playing...  
Radio *g*<sub>a</sub>*g*a(*m*m*a*)!



# Conventional Haloscopes

- 信号

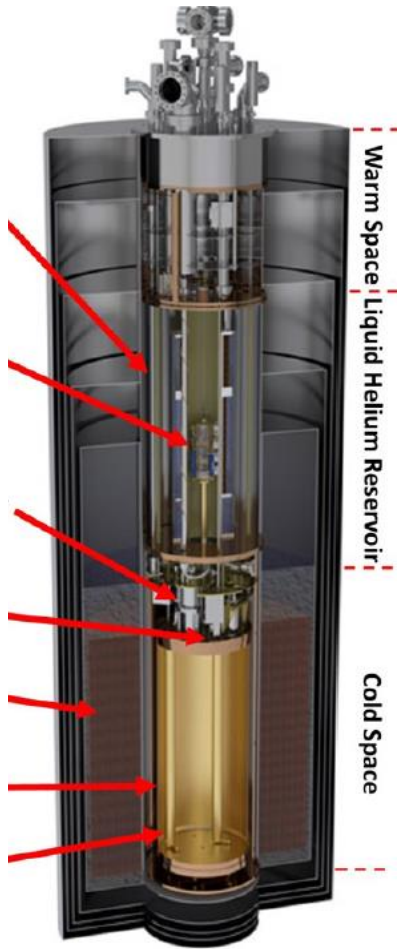
- 共振空洞の利用:  $Q$

$$P_s = 7.2 \times 10^{-23} \text{ W} \left( \frac{\mu\text{eV}}{m_a} \right) \left( \frac{g_{a\gamma}}{2 \times 10^{-16} \text{ GeV}^{-1}} \right)^2 \left( \frac{B_e}{8\text{T}} \right)^2 \left( \frac{\kappa}{0.5} \right) \left( \frac{Q}{10^5} \right) \left( \frac{G_m}{0.69} \right) \left( \frac{V}{200 \text{ l}} \right)$$

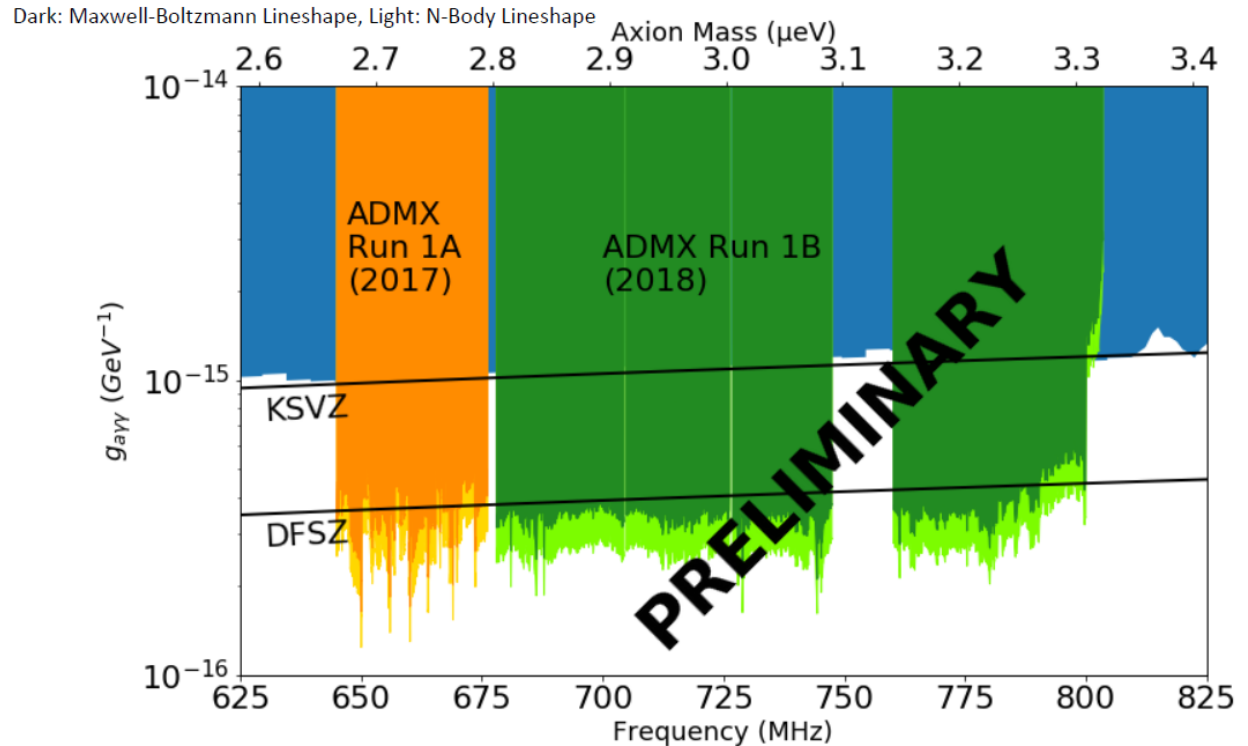
- 実験装置の最適化

- 大きな形状因子 ( $G_m$ ) のためには  $\text{TM}_{010}$  モードが良い
  - 探索質量  $\propto (\text{共振空洞半径})^{-1} \rightarrow V \propto m_a^{-2}$
- 超伝導磁石
  - 強い磁場
  - 大きな内径
- 共振空洞
  - 大きな体積
  - 高いQ値 (ただし上限( $Q_a \approx 10^6$ )あり)

# ADMX



$$V = 136 \text{ l}, B = 6.8 \text{ T}, G_m = 0.4, Q = 50000$$



We estimate sensitivity to DFSZ dark matter axions between 2.8 and 3.3  $\mu\text{eV}$

# Scan rate

- 有限時間内になるべく広い質量領域を探索したい

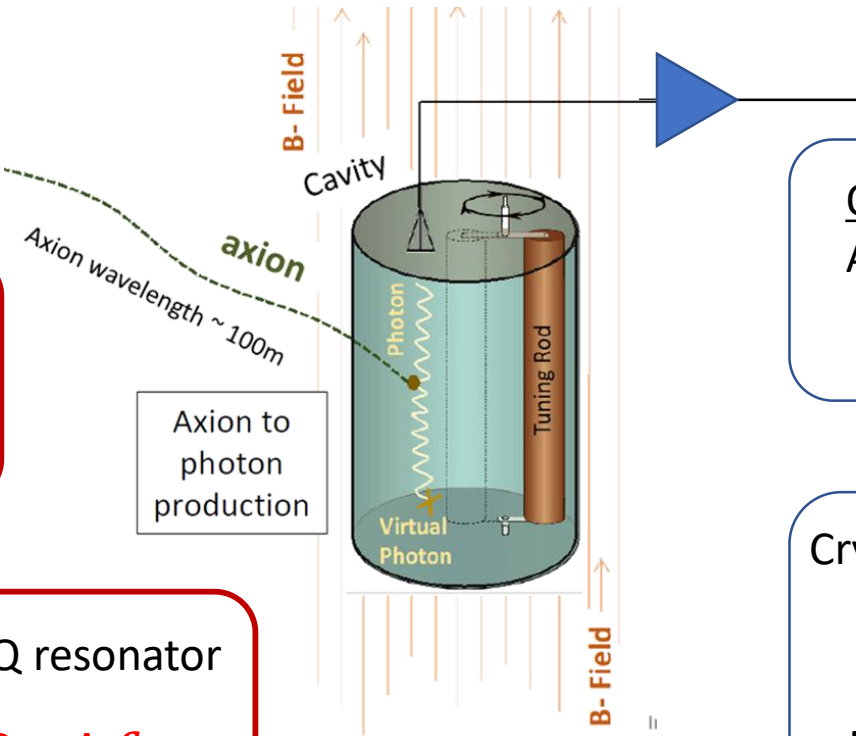


$$\frac{df}{dt} \sim B^4 V^2 G_m^2 Q_L T_{sys}^{-2}$$

$$T_{sys} = T_{phys} + T_{amp}$$

Strong Magnet  
 $B, V$

Tunable High-Q resonator  
 $V, G_m, Q_L, \Delta f$



Quantum-limited Amplifier  
 $T_{amp}$

Cryogenics  
 $T_{phys}$   
Dilution refrigerator

# 戦略

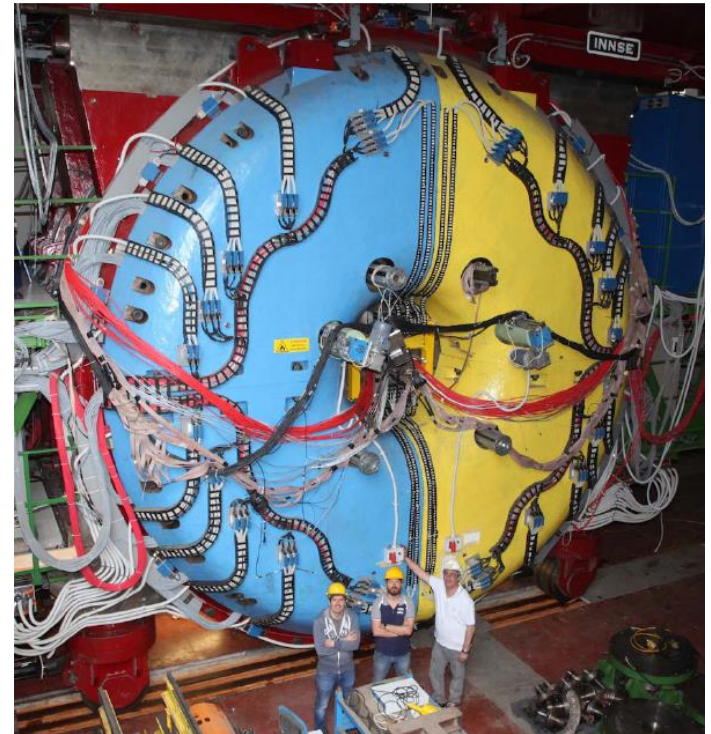
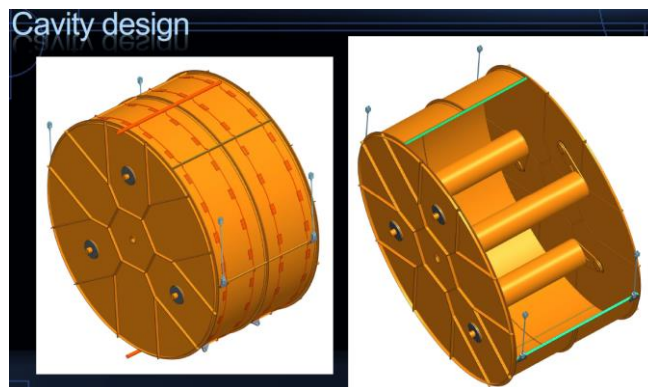
- ADMX... 一番おいしいところ (a few~10  $\mu\text{eV}$ )
  - 空洞サイズ, 磁場, etc.
- 未踏の地は...
  - さらに感度を上げる
  - 軽い方に行くか、重い方に行くか。
  - 軽い方
    - 共振空洞：でかい空洞とボア径大で強いマグネット
    - 他のアイデア：LC回路, Laser干渉計, etc.
  - 重い方
    - 共振空洞
      - 空洞ちっこいので、強いマグネットは作りやすい？
      - でも体積が...複数空洞？
      - 高次のモードを使う...すると形状因子が...
    - 他のアイデア：Dish antenna, Dielectric Haloscope, etc.



# 未探査領域～軽い方～

- でかい共振空洞とでかくて強い磁石！
  - WISPDIMX: DESY HERAの空洞を利用
    - $f_c = 208 \text{ MHz}$ ,  $V \sim 500 \text{ l}$ ,  $Q = 46000$
    - 当面磁場を使わず、hidden photon探索を目指す
  - KLASH: KLOE超伝導磁石の利用
    - $V \sim 22 \text{ m}^3$ ,  $B = 0.6 \text{ T}$

	KLASH	ADMX
$\nu_c$ [MHz]	65-225	450-850
$B^2$ [T <sup>2</sup> ]	0.36	58
$V$ [m <sup>3</sup> ]	22 (5.2)	0.13
$Q_L$	$4 \div 7 \times 10^5$	$0.7 \div 2 \times 10^5$
$\nu_c B^2 V Q_L$ [T <sup>2</sup> m <sup>3</sup> /s]	$0.2 \sim 0.4 \times 10^{15}$	$0.5 \sim 1 \times 10^{15}$



# 未探査領域～さらに軽く～

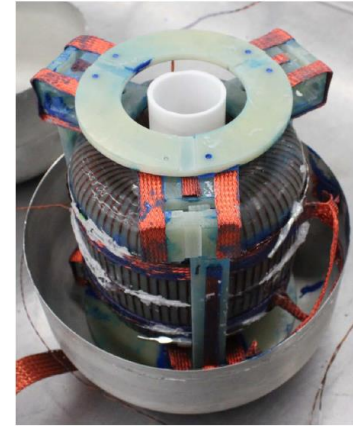
- トロイダル磁石とLC回路

PRL122(2019)121802

- ABRACADABRA

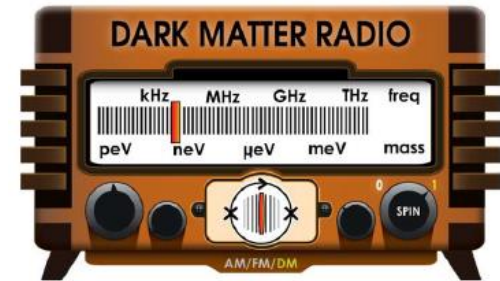
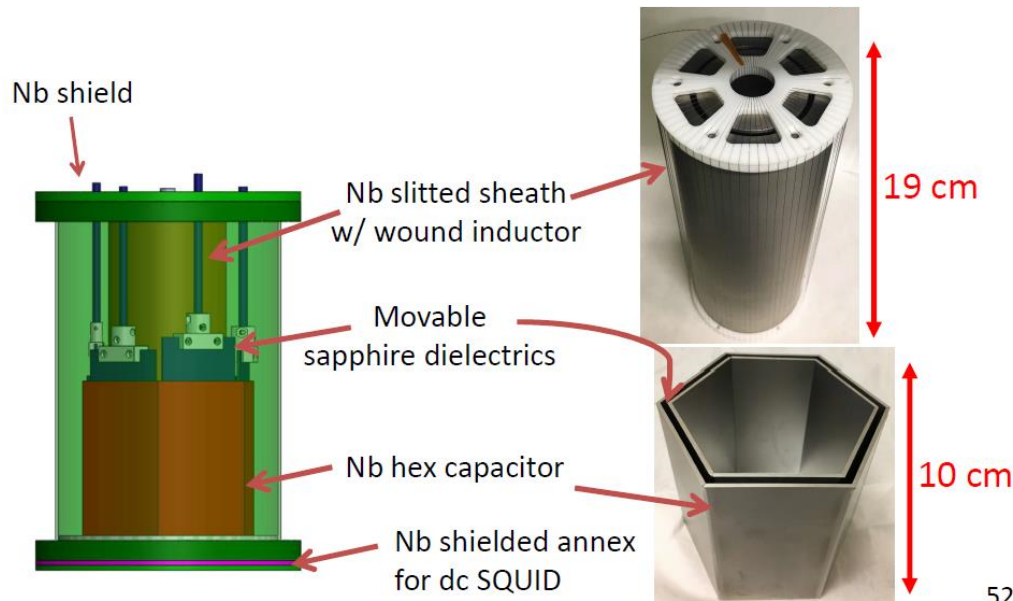


- 最初のステージとしてABRACADABRA-10cm
- まずは、共振(LC)ではなくブロードバンド測定  
質量領域  $3.1 \times 10^{-10} - 8.3 \times 10^{-9}$  eV



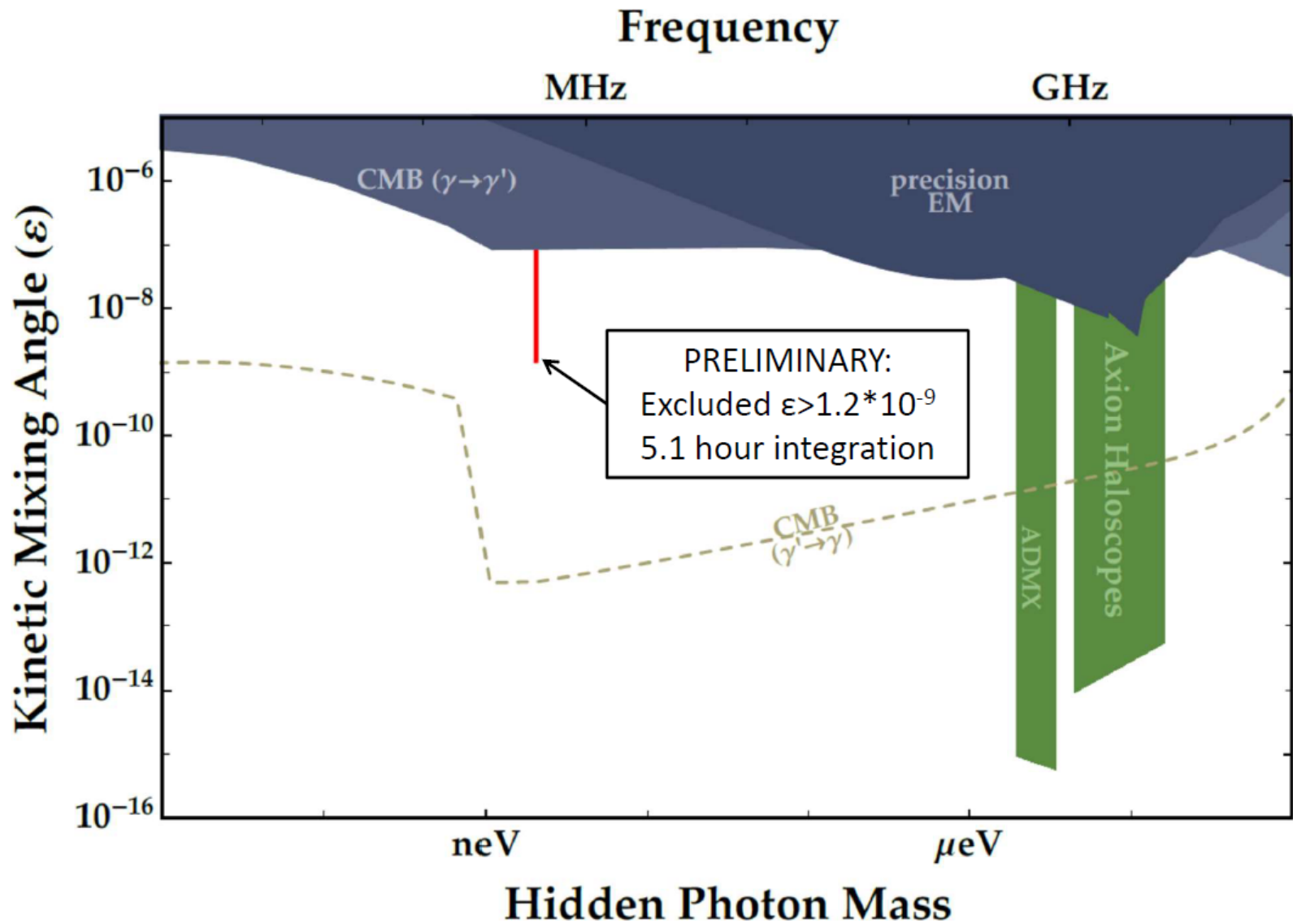
- DM Radio

- 670 mL Pathfinder
- 質量領域：500 peV – 50 neV
- 磁石なしで、hidden photonに焦点



S. Chaudhuri, talk at Caltech HEP Seminar 2019

# DM Radio Pathfinder-preliminary limit!





# 未探査領域～重い方(1)～

## ● 強磁場マグネットの利用

- HAYSTAC

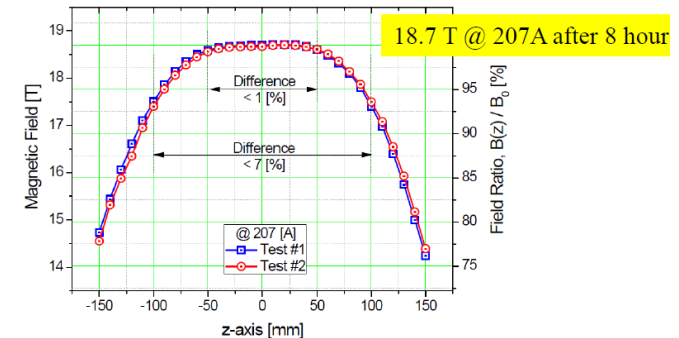
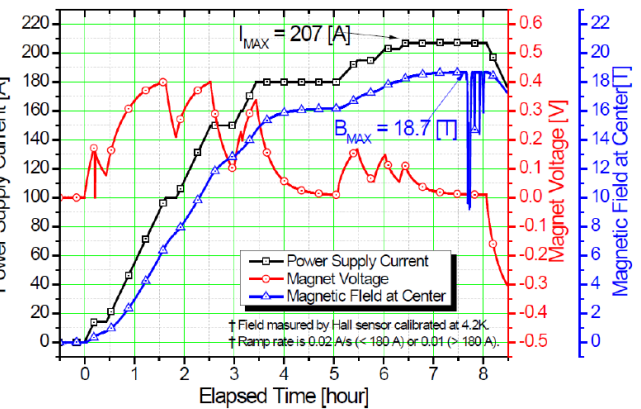
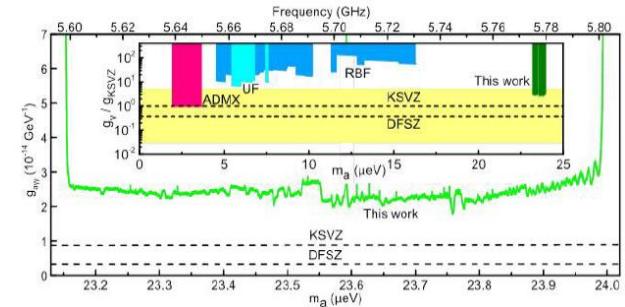
$$f_c = 5.6 - 5.8 \text{ GHz}, V = 1.5 \text{ l}, B = 9 \text{ T}$$

- ORGAN

$$B = 14 \text{ T} \rightarrow 28 \text{ T} \rightarrow \dots$$

- CAPP

$$B = 8 \text{ T} \rightarrow 18 \text{ T} \rightarrow 25 \text{ T} \rightarrow \dots$$

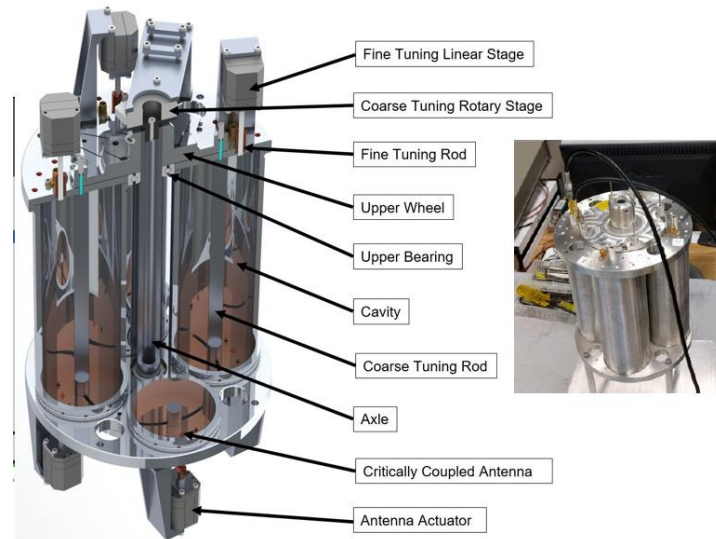


### CAPP Dark Matter Axion Search Schedule

	2017	2018	2019	2020	2021	2022
8T/16cm CAPP8Tb (BlueFos)	setup	experiment		multiple cavity		
18T/7cm CAPP18T (SuNAM)	setup	experiment		multiple cavity		
25T/10cm CAPP25T (BNL)			setup	experiment		
12T/32cm CAPP12T (Oxford)			setup	experiment		

# 未探査領域～重い方(2)～

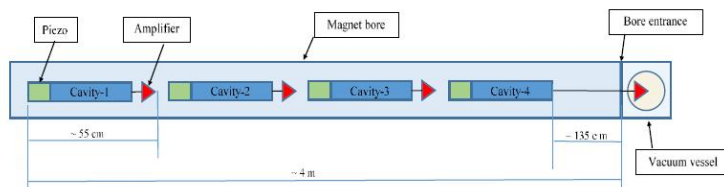
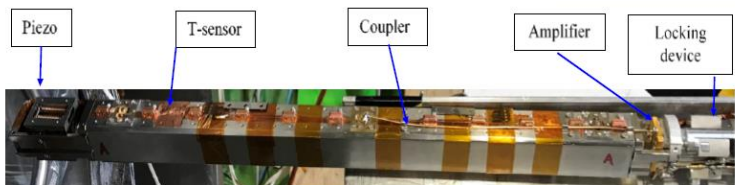
- 既存のマグネット内に複数の共振空洞
  - 空洞間のPhase matchingが必要
  - ADMXのマグネット
    - ADMX-G2で4つの小型空洞
  - CASTのマグネット
    - 長さ(10m)方向に並べる



G. Rybka, talk at PATRAS2019

## CAST-CAPP M. Maroudas, talk at PATRAS2019

4 tunable rectangular stainless steel cavities 25x24x390mm electroplated with  $\sim 30\mu\text{m}$  copper ( $Q_L \sim 10000$ )



## CAST-RADES JMG. Barceló, talk at PATRAS2019

- We have designed and built a concrete implementation of a **30 sub-cavities structure**, that enjoys an optimized resonant mode at  $\sim 8.4\text{ GHz}$ .



# 未探査領域～重い方(3)～

- 共振空洞の高次のモードを使用
  - 形状因子の値が悪くならないように
    - 通常はcancelしてしまう (左)
    - High- $n$  dielectricsの利用 (中)
    - 外部磁場を空間的に変化させる (右)

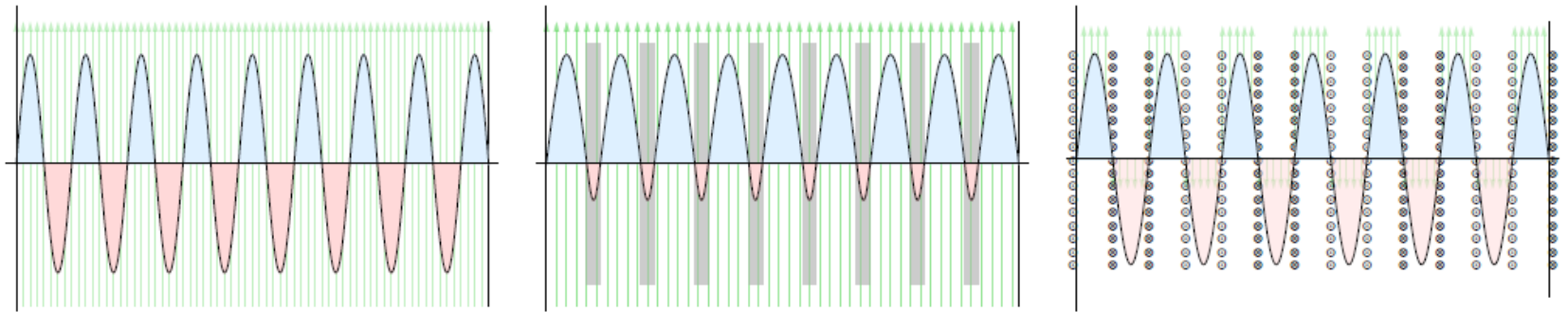
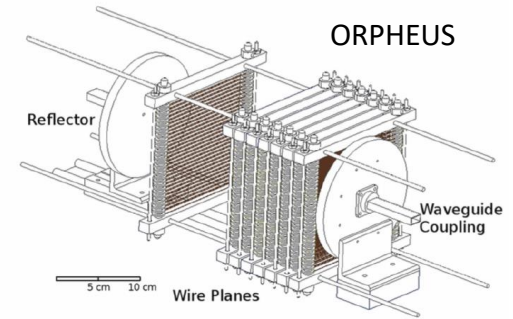
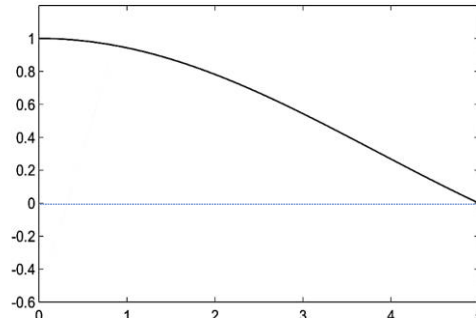
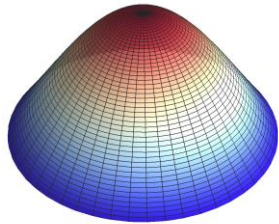


Figure 19: The geometric factor of an ideal 1D cavity in a homogeneous  $B$ -field (green arrows) cancels between crests and valleys of a high mode (left). The cancellation can be avoided by placing high- $n$  dielectrics –grey regions– in the valleys (centre) or by alternating the polarity of the external  $B_e$  field to track the mode variations (right). This case can be done by introducing wire planes with suitable currents [563].

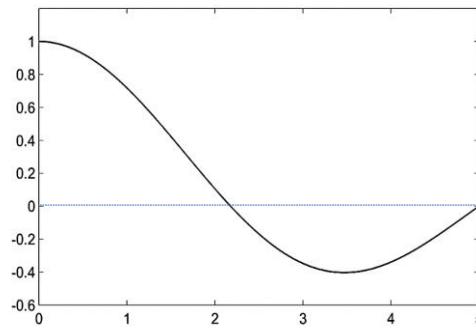
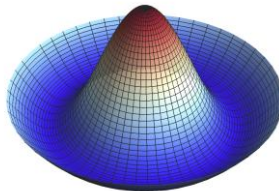
# 高次モードの利用 (CAPP)

SW Youn, talk at PATRAS2019

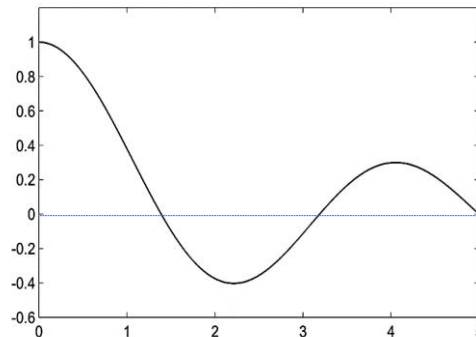
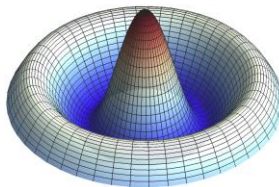
$TM_{010}$



$TM_{020}$



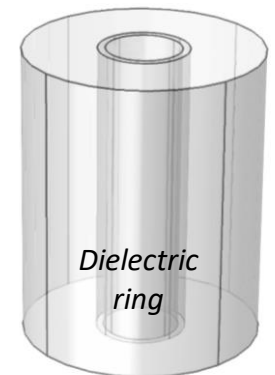
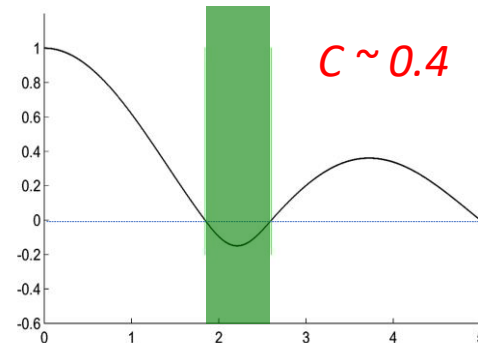
$TM_{030}$



$$P_{a \rightarrow gg} = g_{agg}^2 \frac{r_a}{m_a} B^2 V C_{mnp} \min(Q_L, Q_a)$$

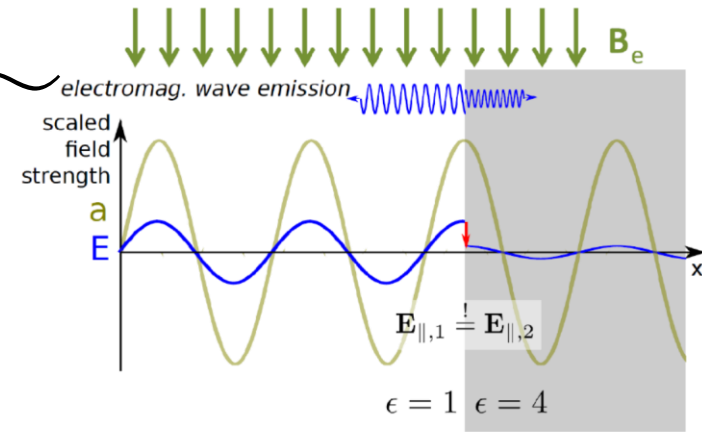
Mode	$f_{rel}$	$Q_{rel}$	$V_{rel}$	$C_{abs}$
$TM_{010}$	1	1	1	0.69
$TM_{020}$	2.3	1.5	1	0.13
$TM_{030}$	3.6	1.9	1	0.05

$$C_{mnp} = \frac{\left| \int \bar{E}_c \cdot \bar{B}_0 dV \right|^2}{\int \epsilon |\bar{E}_c|^2 dV \int |\bar{B}_0|^2 dV}$$



# 未探査領域～重い方(4)～

- 共振空洞はあきらめて...
  - 外部磁場中の境界面( $\epsilon_1 \neq \epsilon_2$ )
    - 光子の放出



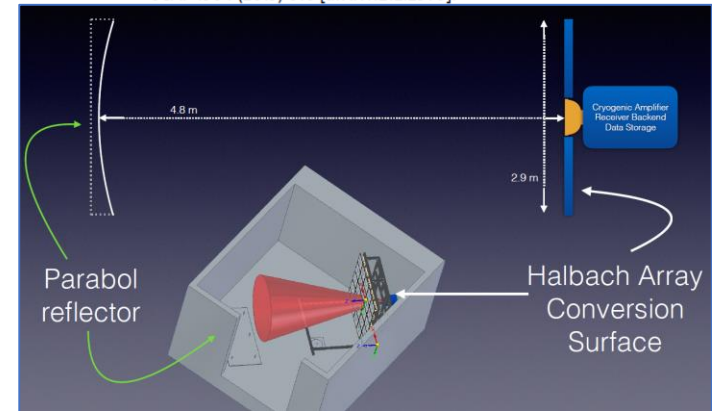
Based on the original idea of:  
 D. Horns, J. Jaeckel, A. Lindner, A. Lobanov, J. Redondo and A. Ringwald  
 JCAP 1304 (2013) 016 [arXiv:1212.2970].

## Dish Antenna

- Broadband search
- BRASS-6実験 (建設中)

## Dielectric Haloscope

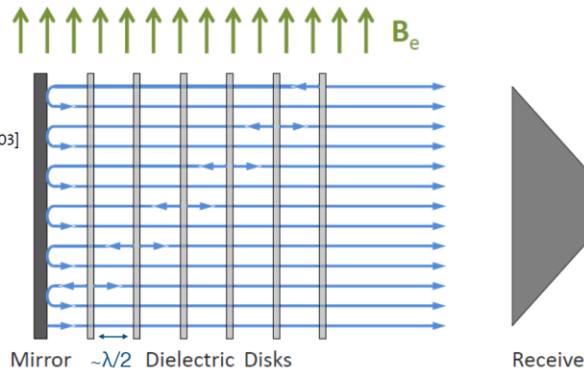
- MADMAX実験 (proof of principle)



Le H. Nguyen, talk at PATRAS2019

- Boost the power by coherent interference of photons generated on N discs plus resonance between discs

Jaeckel and J. Redondo, Phys. Rev. D 88, 115002, (2013) [arXiv:1308.1103]



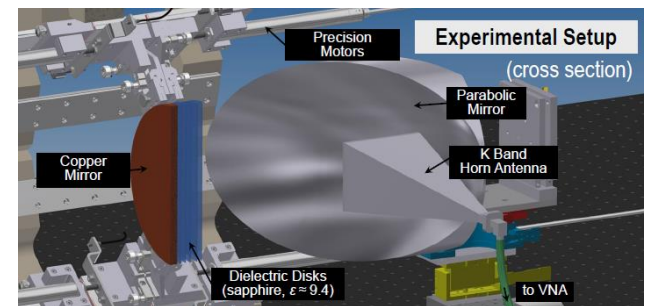
Photon power :

$$\frac{P}{A} = 2 \cdot 10^{-27} \frac{\text{W}}{\text{m}^2} C_{a\gamma\gamma}^2 \left( \frac{B}{10\text{T}} \right)^2 |\beta|^2$$

$$\text{FoM} = B^2 \text{m}^2 = 100 \text{T}^2 \text{m}^2$$

$$\beta^2 = \frac{P_{\text{Diel.Halosc.}}}{P_{\text{Mirror}}}$$

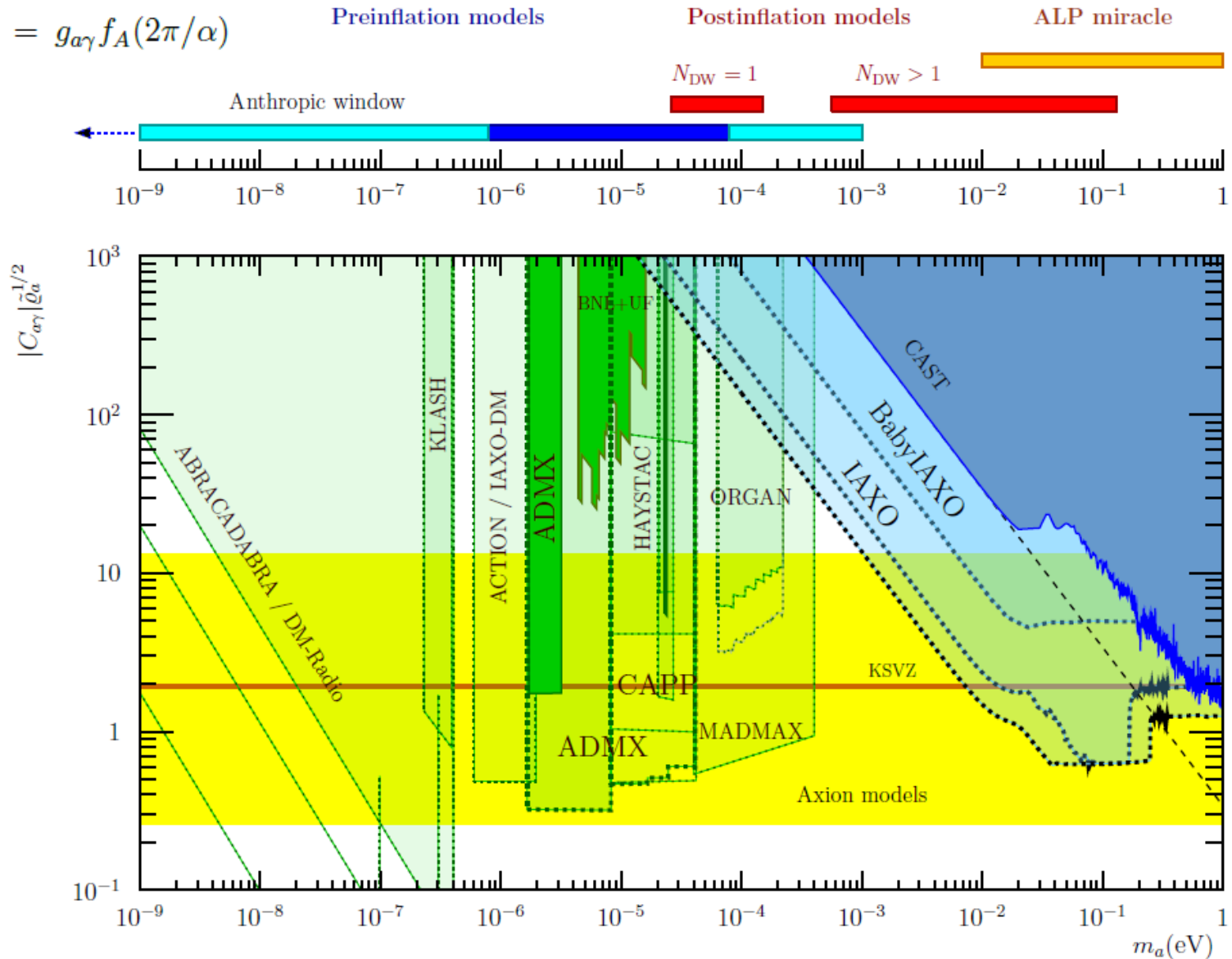
E. Garutti, talk at PATRAS2019



S. Knirck, talk at PATRAS2019

# DH axion探索実験( $g_{a\gamma\gamma}$ )まとめ

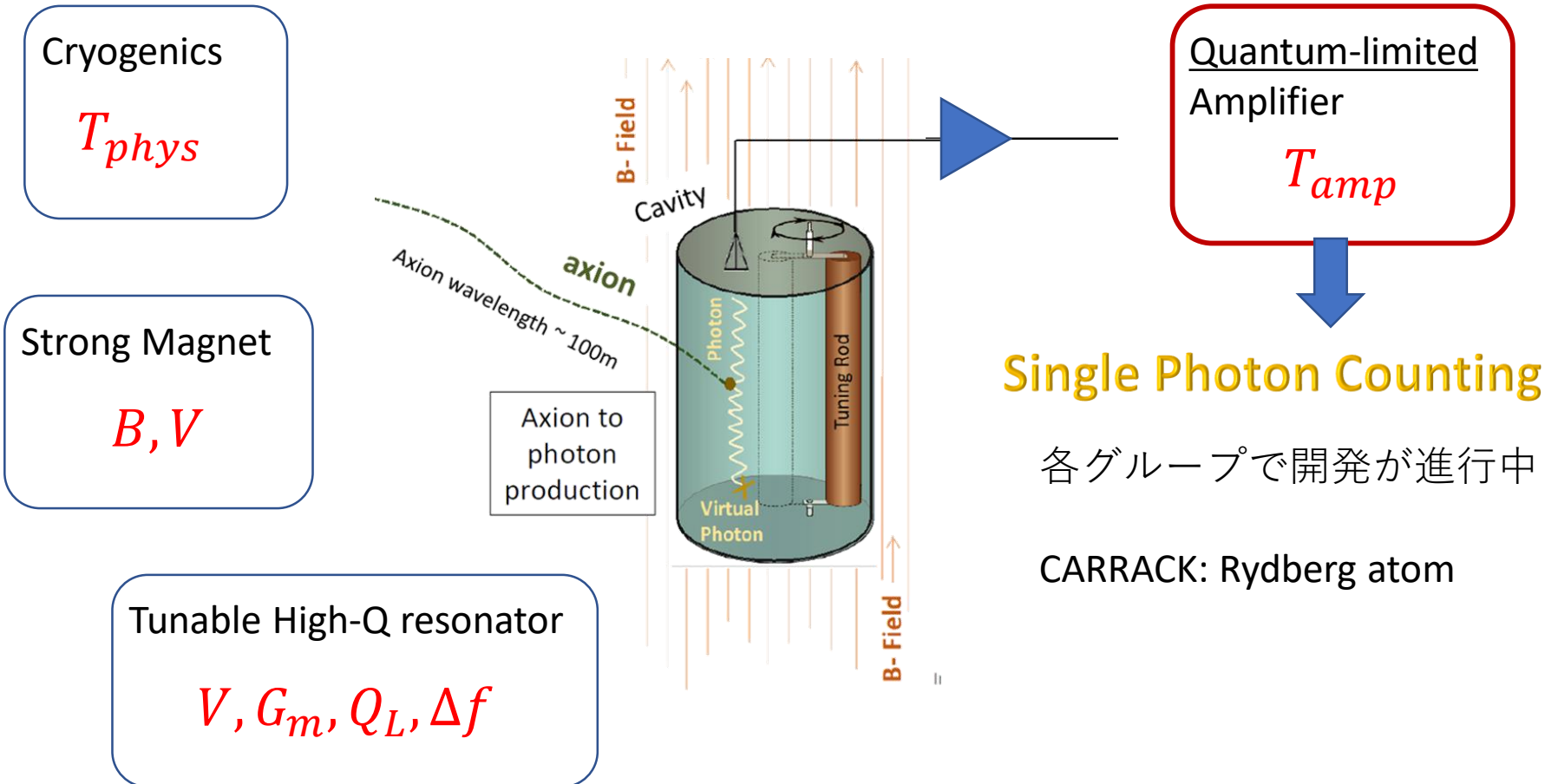
$$C_{a\gamma} = g_{a\gamma} f_A (2\pi/\alpha)$$



# Scan rate

- 有限時間内になるべく広い質量領域を探索したい

$$\frac{df}{dt} \sim B^4 V^2 G_m^2 Q_L T_{sys}^{-2} \quad T_{sys} = T_{phys} + T_{amp}$$



# CARRACK (華洛)

- Rydberg原子によるマイクロ波光子検出
  - SQLなし！
  - Rydberg原子
    - (主としてアルカリ原子で)最外殻電子が主量子数 $n$ の大きい( $\sim 100$ )状態に励起した原子
    - Large radius  $\propto n^2$  ( $R = 0.13\mu m$  for  $n=50$ )
    - Long lifetime  $\tau \propto n^3 \sim \text{msec}$  for  $n \sim 100$
    - Energy level  $E_n \propto n^{-2}$
    - **Energy separation**  $E_{np} - E_{ns} \propto n^{-3}$   
axion as CDM  $\longleftrightarrow$   

$^{39}\text{K}$	$110s_{1/2} \rightarrow 110p_{1/2}$	$10.06 \mu\text{eV}$
$^{85}\text{Rb}$	$112s_{1/2} \rightarrow 112p_{1/2}$	$9.98 \mu\text{eV}$
    - **Large E1 transition probability**  
 $\langle np | e\mathbf{r} | ns \rangle \propto n^2$
    - **Low ionization threshold**  $F_{\text{ionize}} \propto n^{-4}$

classical ionization threshold : 2.2 V/cm for  $n=110$



# Rydberg原子による黒体輻射光子の検出

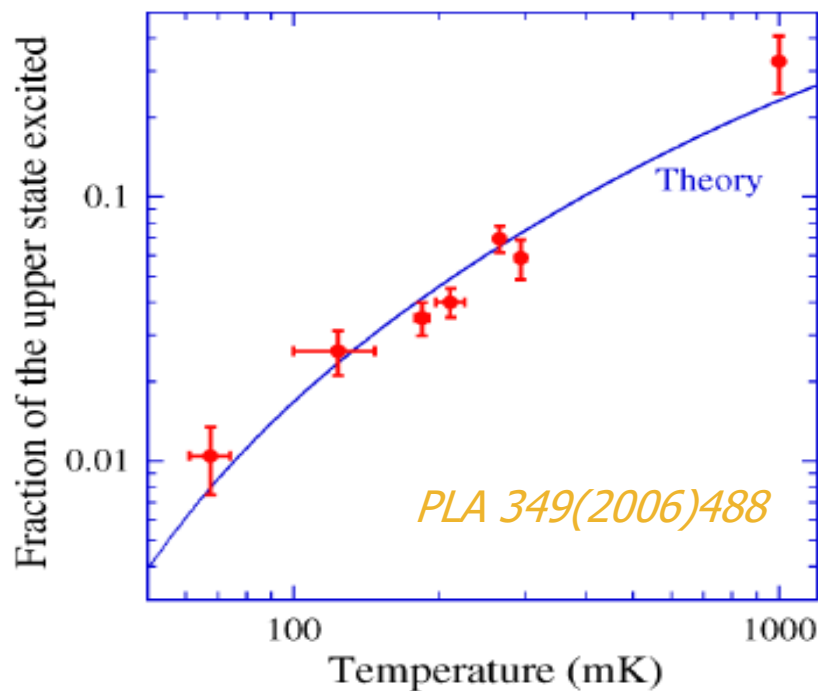
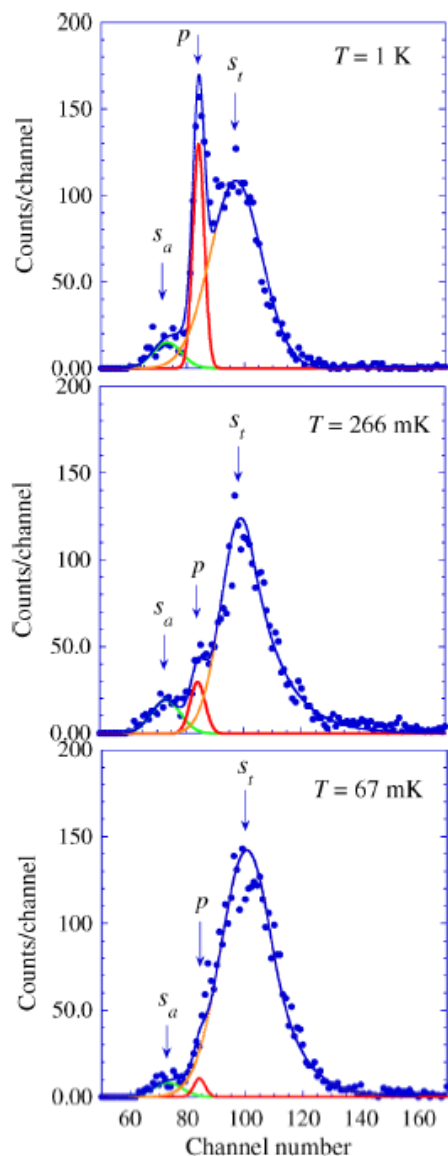
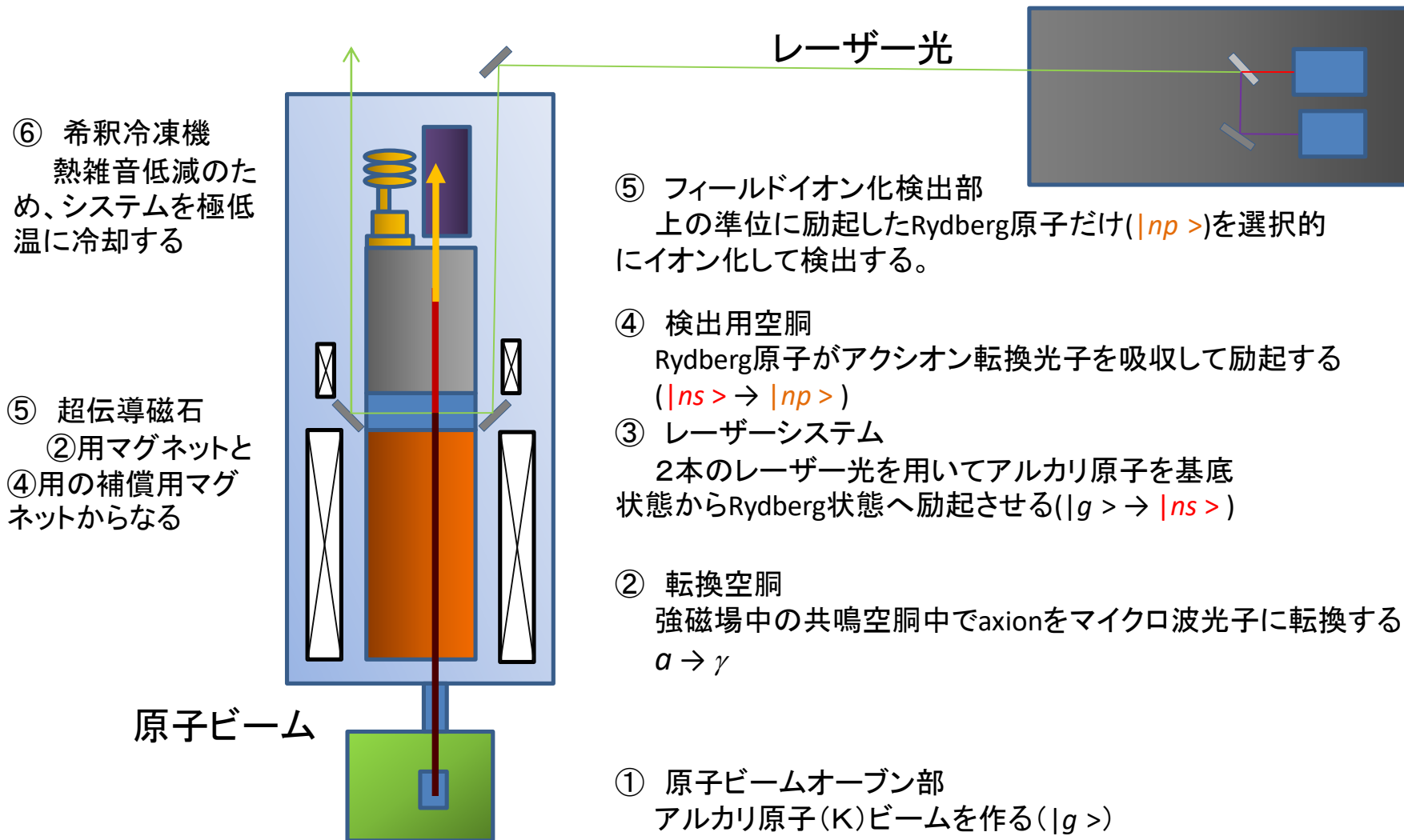


Fig. 4. Temperature dependence of the fraction of the  $111p_{3/2}$  states excited by the blackbody-photon absorption in the cavity. The dependence was measured with the present Rydberg-atom single photon detector. Solid line is a theoretical prediction (see text in detail) in the over-damped regime.

$$111s_{1/2} \Rightarrow 111p_{3/2}$$

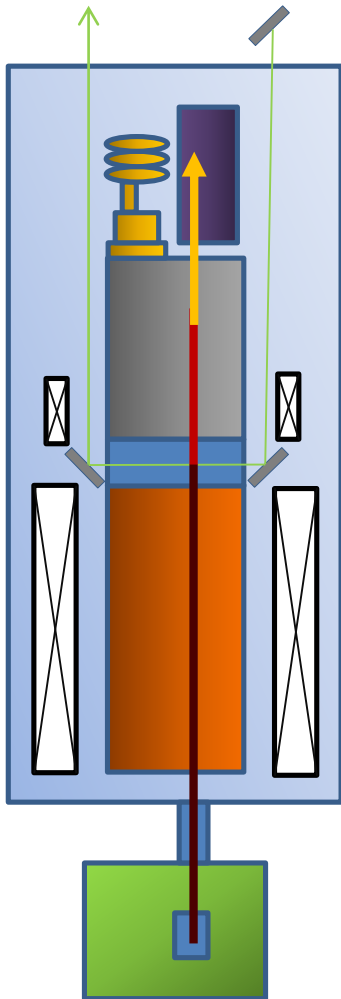
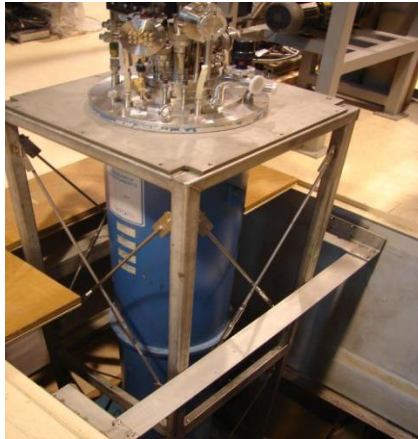
2527 MHz    SQL: 121 mK

# CARRACK装置の原理

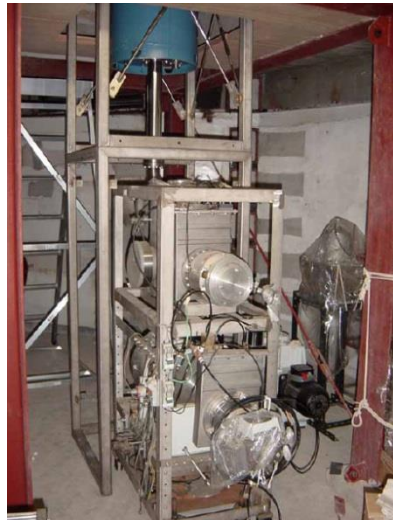


# CARRACK装置

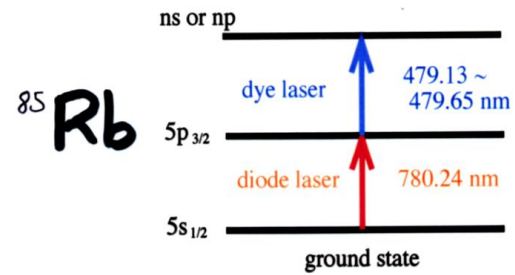
クライオスタット上部



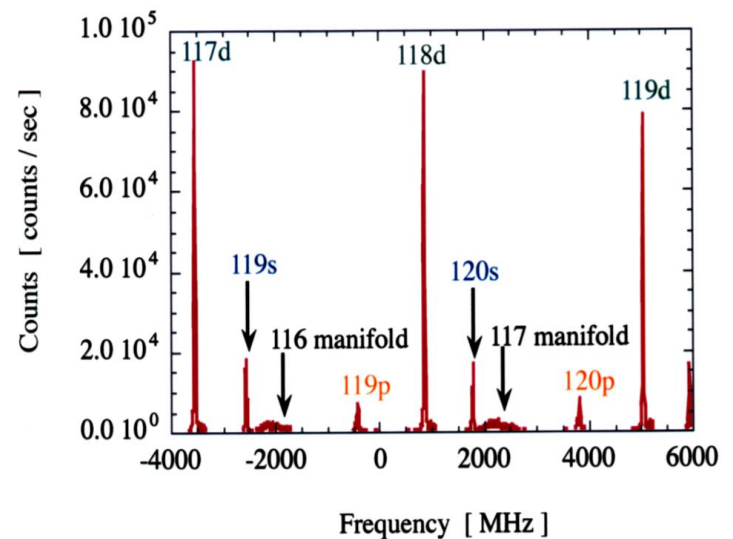
原子ビーム発生装置



Excitation Diagram

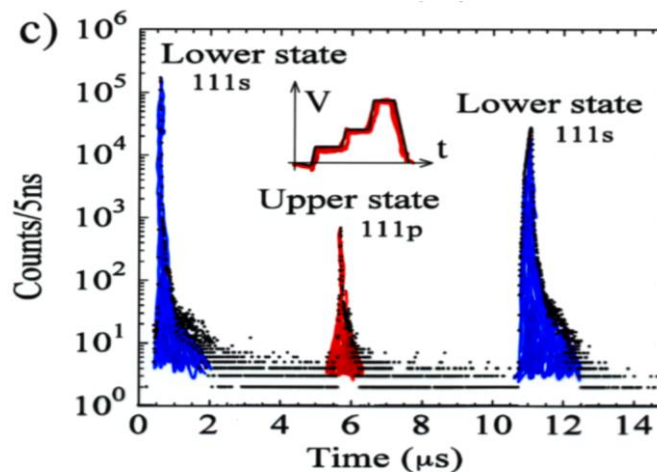
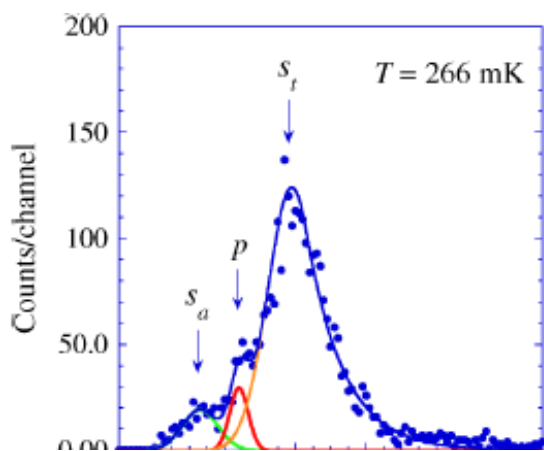


Excitation Spectrum



# Background

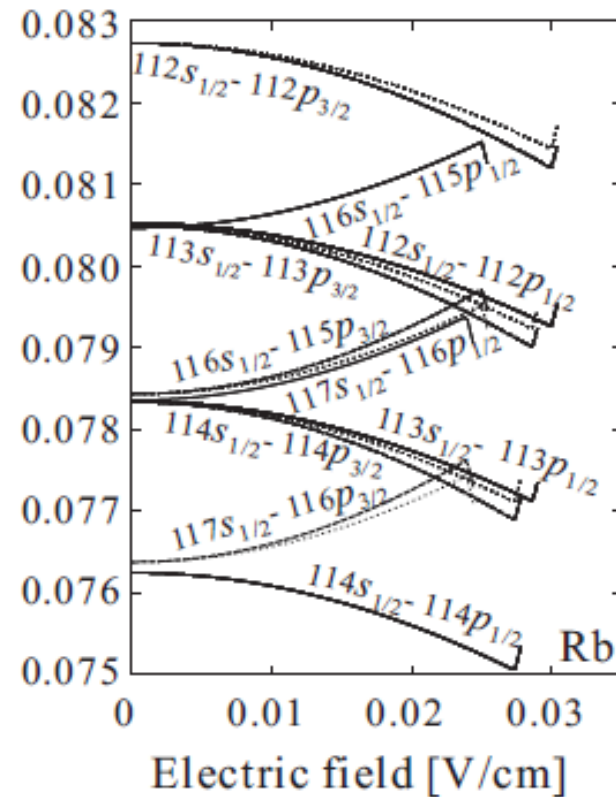
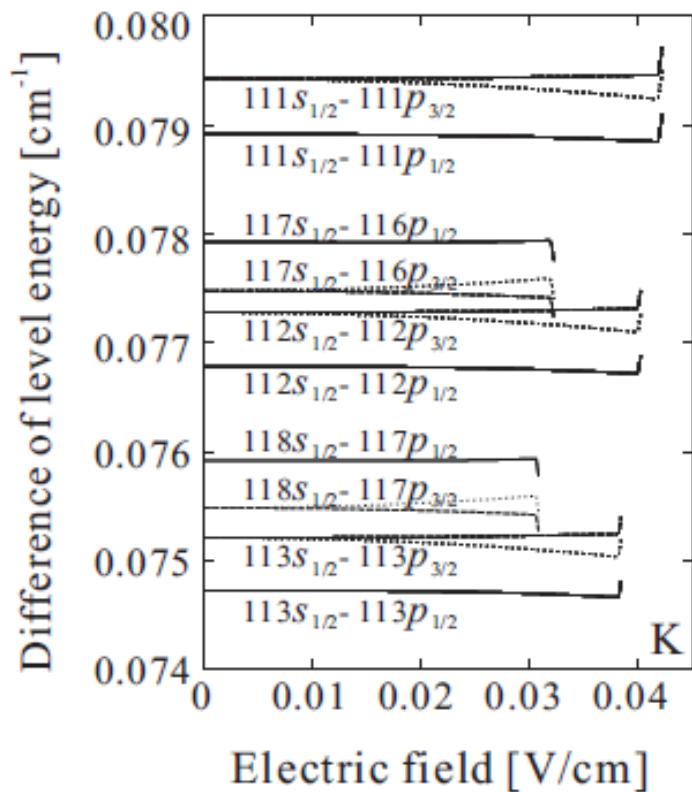
- Cavity中のStray fieldによる準位の混合
    - 極低温時、電極表面に不純物が付着して、空間的に非一様な浮遊電場が生じる。O(10 mV/cm)
- ↓
- 浮遊電場が吸収線幅を広げてしまう。



*Ideal situation*

# newCARRACK

- 原子ビームとして、 $^{85}\text{Rb}$ の代わりに $^{39}\text{K}$ を使用
  - 外部電場に対して非常に鈍感であると予想されている



# 原子ビームのバンチ化

- *CARRACKの改良点*

- 原子をRbから外部電場に鈍感な $^{39}\text{K}$ に変更
- 探索質量域の変更 (2.4 GHz  $\rightarrow$  24 GHz)
  - これにともなう**検出効率向上の必要性**



- *バンチ化による検出効率の向上*

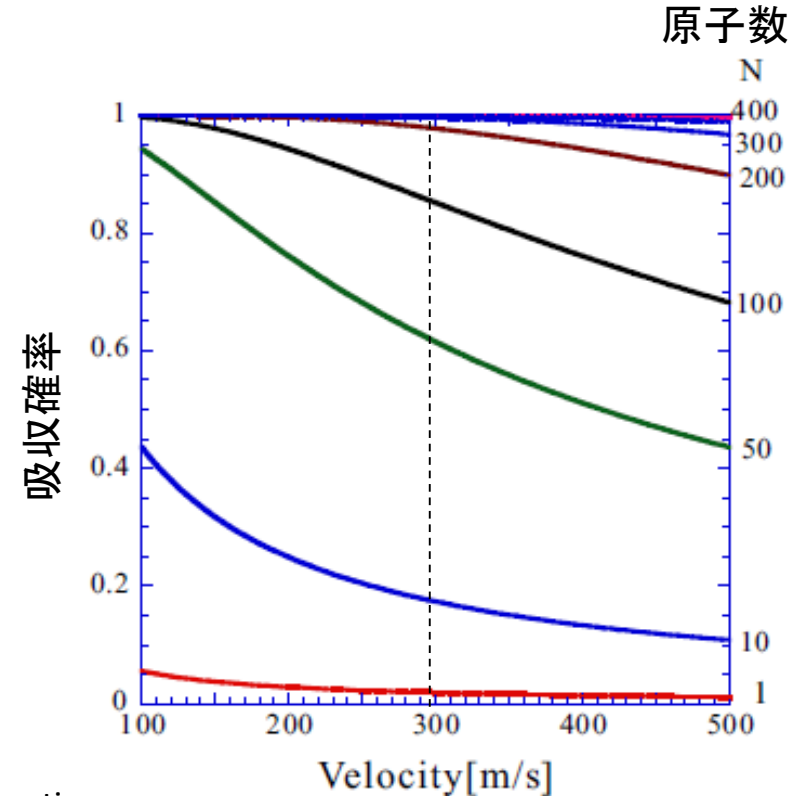
- フィールドイオン化パルス電場にタイミングを合わせたバンチを供給  $\rightarrow$  duty cycleを上げる。
- 転換光子を100%吸収できる原子数が必要。



原子ビームを高輝度化した上で、バンチ化する

# アクション検出に必要な原子ビーム

- 100 %の吸収確率
  - 1バンチ内の原子数 : >200 atoms/bunch
- Rydberg状態の維持(寿命: ~1 ms)
  - ビーム速度: 300 m/s
- 幾何学的問題(~3 mの飛行)
  - 速度広がり: 1 m/s(縦横ともに)

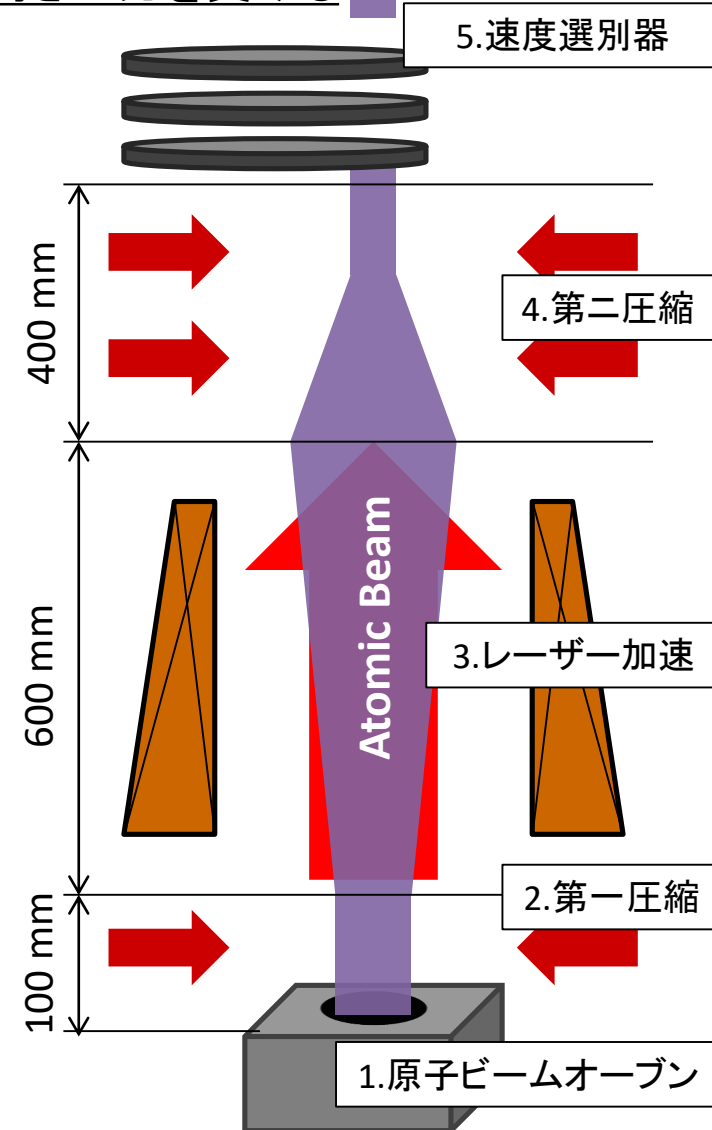


# バンチ化原子ビーム生成方法

原子ビームはレーザー(赤矢印)の向きに力を受ける

## バンチ化原子ビーム生成プロセス

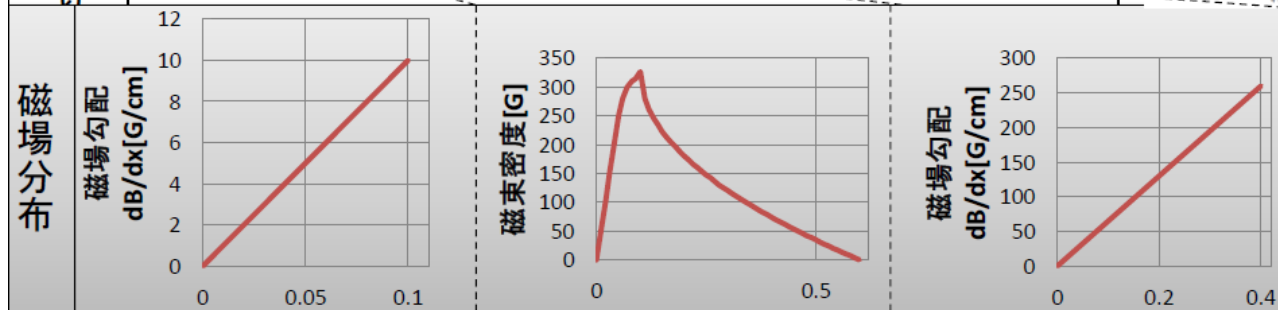
1. 原子ビームオープン  
→ 原子ビーム発生
2. レーザー冷却(第一圧縮)  
→ レーザー加速の間の広がりを抑えるために、横方向の運動量を圧縮。
3. レーザー加速  
→ 速度の遅い原子を加速して300 m/sにする。
4. レーザー冷却(第二圧縮)  
→ 原子ビーム高輝度化。
5. 速度選別器(ディスクスロット型)  
→ 速度300m/sの原子のみを選別し、バンチ化



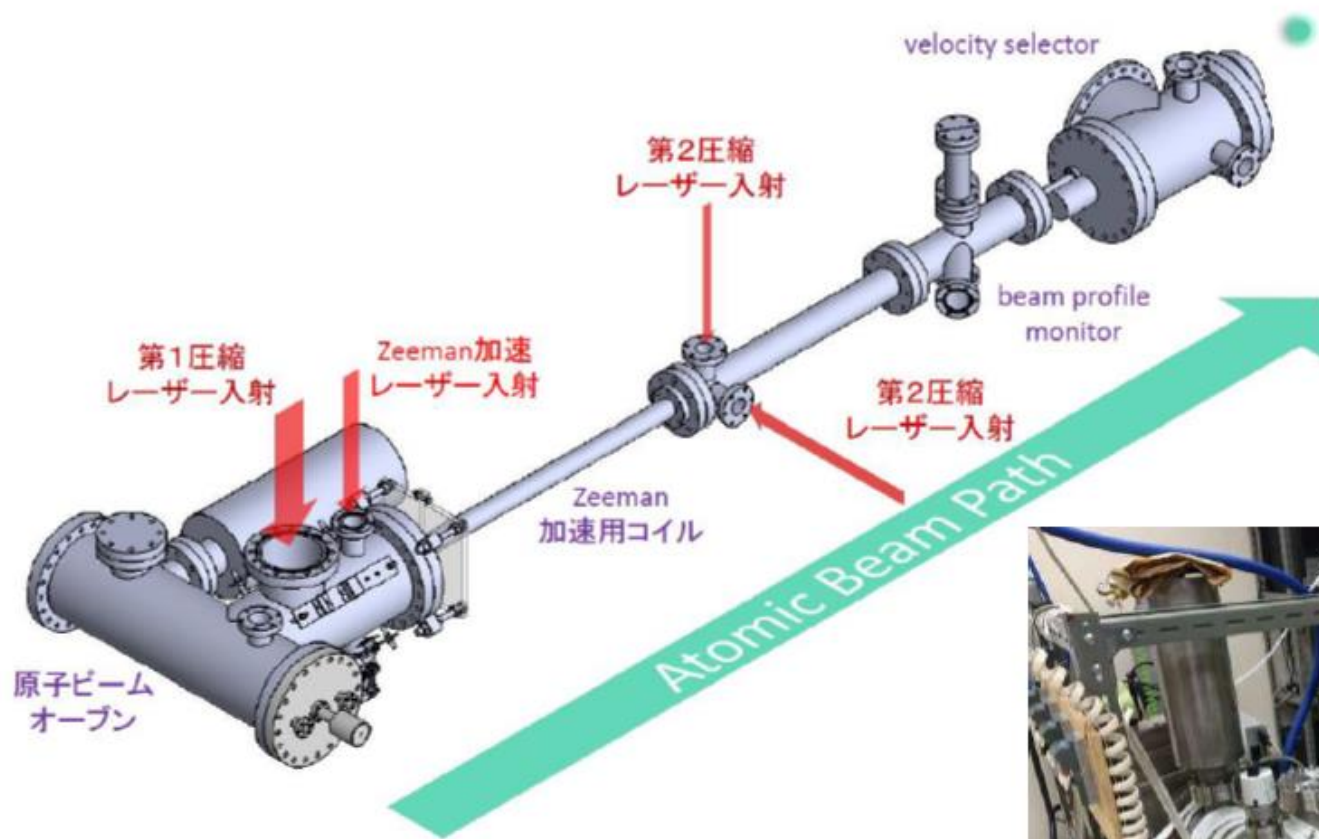


# 原子ビームシミュレーション

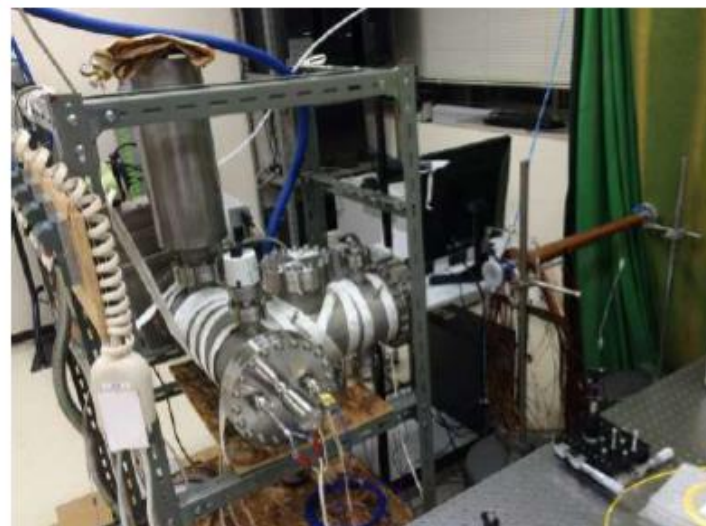
原子ビームシミュレーション		分布図	計算 [mm]	要請 [mm]
横位置広がり			5	4
横速度広がり			0.8	1
縦速度広がり			1	1



# Rydberg原子ビームのバンチ化

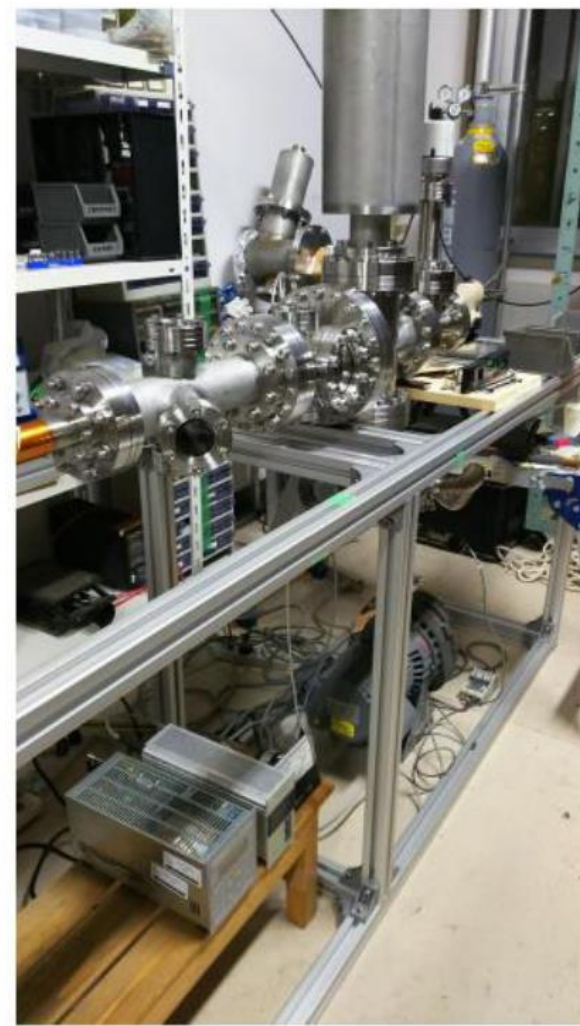


- 機器の一部は製作済み。
- テストはまだ。



# バンチ化ビーム装置

- 京都大学にて建設中
- レーザー系と合わせてテストを行う
  - ビームプロファイルモニターで性能確認予定



# まとめ

- アンプ系の雑音を下げる
  - 単光子検出が有効
    - 世界中で開発・テスト中
  - Rydberg原子は単光子検出が可能
  - 検出効率を向上すべくK原子のバンチ化ビーム発生装置を京大にて構築中